

A mixed finite element method based on a biorthogonal system for nearly incompressible elastic problems

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Abstract

A Petrov–Galerkin scheme in a saddle point formulation gives rise to a non-symmetric saddle point problem. This article considers a non-symmetric saddle point problem with a penalty parameter. A mixed finite element method for linear elasticity based on a Petrov–Galerkin formulation is then analyzed within the framework of the non-symmetric saddle point problem with penalty. Working with a biorthogonal system to discretize the pressure equation, we obtain a robust and efficient numerical scheme for nearly incompressible linear elasticity using linear finite elements. A numerical example demonstrates the robustness of the approach. These results are useful to analyze a Petrov–Galerkin scheme in a saddle point problem.

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1 Introduction and abstract setting

Low order finite elements based on standard displacement formulation exhibit a locking effect when applied to nearly incompressible problems [1, 2]. A standard remedy for the locking effect is to work with mixed methods introducing some extra variables leading to a saddle point problem [3, 4]. Here we consider a simple mixed formulation in displacement–pressure form [4], and apply a Petrov–Galerkin formulation for the pressure equation. The pressure equation is discretized using different trial and test spaces, where the bases of the trial and test spaces form a biorthogonal system. This results in a diagonal mass matrix for the pressure equation, and it is easy to statically condense out the pressure inverting a diagonal matrix. We proposed this formulation in recent work [5] as an alternative approach to the method based on primal and dual meshes. However, the mathematical analysis presented in that study covers only the case of primal and dual meshes. We note that the method based on primal and dual meshes results in a symmetric saddle point system, whereas the Petrov–Galerkin formulation for the pressure equation yields a non-symmetric saddle point system. The novel idea of this article is to show the well-posedness of this formulation by combining the results on the non-symmetric saddle point problem [6, 7, 8] and the stability result of mini-element [9].

We start with an abstract setting. Let V , W , P and Q be Hilbert spaces with inner products $(\cdot, \cdot)_V$, $(\cdot, \cdot)_W$, $(\cdot, \cdot)_P$ and $(\cdot, \cdot)_Q$, respectively. Let $\mathbf{a}(\cdot, \cdot) : V \times W \rightarrow \mathbb{R}$, $\mathbf{b}_1(\cdot, \cdot) : W \times P \rightarrow \mathbb{R}$, $\mathbf{b}_2(\cdot, \cdot) : V \times Q \rightarrow \mathbb{R}$, and $\mathbf{c}(\cdot, \cdot) : P \times Q \rightarrow \mathbb{R}$ be bilinear forms. We consider a non-symmetric saddle point problem with penalty: given $f \in W'$ and $g \in Q'$, find $(\mathbf{u}, \mathbf{p}) \in V \times P$ so that

$$\begin{aligned} \mathbf{a}(\mathbf{u}, \mathbf{w}) + \mathbf{b}_1(\mathbf{w}, \mathbf{p}) &= f(\mathbf{w}), \quad \mathbf{w} \in W, \\ \mathbf{b}_2(\mathbf{u}, \mathbf{q}) - \mathbf{t} \mathbf{c}(\mathbf{p}, \mathbf{q}) &= g(\mathbf{q}), \quad \mathbf{q} \in Q, \end{aligned} \quad (1)$$

where \mathbf{t} is a positive small parameter, and W' and Q' denote the space of continuous linear functionals on W and Q , respectively. We are interested in analyzing the well-posedness of the problem (1) when $\mathbf{t} \rightarrow 0$. There are a number of publications devoted to the analysis of this problem in different particular forms; for example, when $\mathbf{b}_1(\cdot, \cdot)$ and $\mathbf{b}_2(\cdot, \cdot)$ are the same, and $V = W$ and $P = Q$ [3, 4]. The case with $\mathbf{t} = 0$ was analyzed by Nicolaidis [6] and Bernardi et al. [7], whereas the case with $V = W$, $P = Q$ and $\mathbf{t} = 1$ was considered by Ciarlet et al. [8]. We show the stability of the problem (1) by combining the ideas presented by Ciarlet et al. [8] and Bernardi et al. [7]. To this end, we assume that the bilinear forms $\mathbf{a}(\cdot, \cdot)$, $\mathbf{b}_1(\cdot, \cdot)$, $\mathbf{b}_2(\cdot, \cdot)$ and $\mathbf{c}(\cdot, \cdot)$ satisfy

$$\begin{aligned} |\mathbf{a}(\mathbf{v}, \mathbf{w})| &\leq \bar{\mathbf{a}} \|\mathbf{v}\|_V \|\mathbf{w}\|_W, \quad \mathbf{v} \in V, \mathbf{w} \in W, \\ |\mathbf{b}_1(\mathbf{w}, \mathbf{p})| &\leq \bar{\mathbf{b}}_1 \|\mathbf{w}\|_W \|\mathbf{p}\|_P, \quad \mathbf{w} \in W, \mathbf{p} \in P, \\ |\mathbf{b}_2(\mathbf{v}, \mathbf{q})| &\leq \bar{\mathbf{b}}_2 \|\mathbf{v}\|_V \|\mathbf{q}\|_Q, \quad \mathbf{v} \in V, \mathbf{q} \in Q, \\ |\mathbf{c}(\mathbf{p}, \mathbf{q})| &\leq \bar{\mathbf{c}} \|\mathbf{p}\|_P \|\mathbf{q}\|_Q, \quad \mathbf{p} \in P, \mathbf{q} \in Q, \end{aligned} \quad (2)$$

where $\bar{\mathbf{a}}$, $\bar{\mathbf{b}}_1$, $\bar{\mathbf{b}}_2$ and $\bar{\mathbf{c}}$ are continuity constants of the bilinear forms $\mathbf{a}(\cdot, \cdot)$, $\mathbf{b}_1(\cdot, \cdot)$, $\mathbf{b}_2(\cdot, \cdot)$ and $\mathbf{c}(\cdot, \cdot)$, respectively, and $\|\cdot\|_V$, $\|\cdot\|_W$, $\|\cdot\|_P$ and $\|\cdot\|_Q$ are induced inner products from the associated norms.

We define $\mathbf{U}_W \subset W$ and $\mathbf{U}_V \subset V$ as

$$\begin{aligned} \mathbf{U}_W &:= \{\mathbf{w} \in W : \mathbf{b}_1(\mathbf{w}, \mathbf{p}) = 0, \mathbf{p} \in P\}, \\ \mathbf{U}_V &:= \{\mathbf{v} \in V : \mathbf{b}_2(\mathbf{v}, \mathbf{q}) = 0, \mathbf{q} \in Q\}, \end{aligned}$$

and assume that

$$\begin{aligned}
 \sup_{w \in U_W} \frac{a(v, w)}{\|w\|_W} &\geq \alpha \|v\|_V, \quad v \in U_V, \\
 \sup_{v \in U_V} a(v, w) &> 0, \quad w \in U_W \setminus \{0\}, \\
 \sup_{w \in W} \frac{b_1(w, p)}{\|w\|_W} &\geq \beta_1 \|p\|_P, \quad p \in P, \\
 \sup_{v \in V} \frac{b_2(v, q)}{\|v\|_V} &\geq \beta_2 \|q\|_Q, \quad q \in Q
 \end{aligned} \tag{3}$$

hold for some constants $\alpha, \beta_1, \beta_2 > 0$, where the supremum is taken only over the non-trivial elements of the underlying sets. In order to show the existence and uniqueness of the problem (1), we need the following theorem, which was proved by Nicolaides [6] and Bernardi et al. [7].

Theorem 1 *Let assumptions (2) and (3) be satisfied. Then for any $f \in W'$ and $g \in Q'$, there exists a unique solution $(u, p) \in V \times P$ to the saddle point problem of finding $(u, p) \in V \times P$ so that*

$$\begin{aligned}
 a(u, w) + b_1(w, p) &= f(w), \quad w \in W, \\
 b_2(u, q) &= g(q), \quad q \in Q,
 \end{aligned} \tag{4}$$

which satisfies the following stability estimates:

$$\|u\|_V \leq \beta_2^{-1} (1 + \alpha^{-1} \bar{a}) \|g\|_{Q'} + \alpha^{-1} \|f\|_{W'}, \quad \|p\|_P \leq \beta_1^{-1} (\|f\|_{W'} + \bar{a} \|u\|_V). \tag{5}$$

The existence and uniqueness of the problem (1) is then established by the following theorem. Although we use a non-symmetric formulation in contrast to Ciarlet et al. [8], the proof of this theorem is similar [8, Theorem 3.2].

Theorem 2 *Let assumptions (2) and (3) be satisfied, and*

$$\delta := \beta_1^{-1} \beta_2^{-1} \bar{a} (1 + \alpha^{-1} \bar{a}) t \bar{c} < 1. \quad (6)$$

Then for any $f \in V'$ and $g \in Q'$, there exists a unique solution $(\mathbf{u}, \mathbf{p}) \in V \times P$ to the saddle point problem (1) satisfying the following stability estimates:

$$\|\mathbf{p}\|_P \leq \frac{1}{1 - \delta} \|\tilde{\mathbf{p}}\|_P, \quad \|\mathbf{u}\|_V \leq \|\tilde{\mathbf{u}}\|_V + \frac{\beta_2 (1 + \alpha^{-1} \bar{a}) t \bar{c}}{1 - \delta} \|\tilde{\mathbf{p}}\|_P, \quad (7)$$

where $(\tilde{\mathbf{u}}, \tilde{\mathbf{p}})$ is the solution to (4) and satisfies the bounds

$$\|\tilde{\mathbf{u}}\|_V \leq \beta_2^{-1} (1 + \alpha^{-1} \bar{a}) \|g\|_{Q'} + \alpha^{-1} \|f\|_{W'}, \quad \|\tilde{\mathbf{p}}\|_P \leq \beta_1^{-1} (\|f\|_{W'} + \bar{a} \|\tilde{\mathbf{u}}\|_V).$$

Proof: Letting $\mathbf{p}_0 = \mathbf{0} \in P$, we define a sequence $\{(\mathbf{u}_n, \mathbf{p}_n)\}$ for $n \in \mathbb{N}$ by

$$\begin{aligned} \mathbf{a}(\mathbf{u}_{n+1}, \mathbf{w}) + \mathbf{b}_1(\mathbf{w}, \mathbf{p}_{n+1}) &= f(\mathbf{w}), \quad \mathbf{w} \in W \\ \mathbf{b}_2(\mathbf{u}_{n+1}, \mathbf{q}) &= g(\mathbf{q}) + t \mathbf{c}(\mathbf{p}_n, \mathbf{q}), \quad \mathbf{q} \in Q. \end{aligned} \quad (8)$$

The sequence is well-defined from Theorem 1, and for $n \in \mathbb{N}$ we have

$$\begin{aligned} \mathbf{a}(\mathbf{u}_{n+1} - \mathbf{u}_n, \mathbf{w}) + \mathbf{b}_1(\mathbf{w}, \mathbf{p}_{n+1} - \mathbf{p}_n) &= 0, \quad \mathbf{w} \in W \\ \mathbf{b}_2(\mathbf{u}_{n+1} - \mathbf{u}_n, \mathbf{q}) &= t \mathbf{c}(\mathbf{p}_n - \mathbf{p}_{n-1}, \mathbf{q}), \quad \mathbf{q} \in Q. \end{aligned} \quad (9)$$

Theorem 1 yields the existence and uniqueness of the solution of (9) with the estimates

$$\begin{aligned} \|\mathbf{u}_{n+1} - \mathbf{u}_n\|_V &\leq \beta_2^{-1} (1 + \alpha^{-1} \bar{a}) t \bar{c} \|\mathbf{p}_n - \mathbf{p}_{n-1}\|_P, \\ \|\mathbf{p}_{n+1} - \mathbf{p}_n\|_P &\leq \beta_1^{-1} \bar{a} \|\mathbf{u}_{n+1} - \mathbf{u}_n\|_V, \end{aligned} \quad (10)$$

and hence

$$\|\mathbf{p}_{n+1} - \mathbf{p}_n\|_P \leq \beta_1^{-1} \beta_2^{-1} \bar{a} (1 + \alpha^{-1} \bar{a}) t \bar{c} \|\mathbf{p}_n - \mathbf{p}_{n-1}\|_P \leq \delta^n \|\mathbf{p}_1\|_P. \quad (11)$$

Now taking $n \in \mathbb{N}$ and an integer $m > n$, we have

$$\|p_m - p_n\|_P \leq \sum_{i=n}^{m-1} \|p_{i+1} - p_i\|_P \leq \sum_{i=n}^{m-1} \delta^i \|p_1\|_P \leq \frac{\delta^n}{1-\delta} \|p_1\|_P, \quad (12)$$

which shows that $\{p_n\}$ is a Cauchy sequence, and so converges to a $p \in P$. The stability estimate for p is obtained by taking $n = 0$ in (12). Using the first inequality of (10) and the estimate (11), the sequence $\{u_n\}$ is shown to be a Cauchy sequence, and stability estimate for u is obtained similarly as for p . The uniqueness and other details is to be worked out as in the work of Ciarlet et al. [8, Theorem 3.2]. ♠

Our primary concern is to find a robust approximation scheme based on linear finite elements and simplicial triangulation for the nearly incompressible elastic problem. In particular, we are interested in a mixed scheme where the pressure variable is eliminated and obtain a formulation based only on the displacement variable.

Let $\Omega \subset \mathbb{R}^d$ with $d \in \{2, 3\}$, $L^2(\Omega)$ be the space of square-integrable functions defined on Ω with the inner product and norm being denoted by $(\cdot, \cdot)_0$ and $\|\cdot\|_0$, respectively, and $L^2_0(\Omega) := \{p \in L^2(\Omega) : \int_{\Omega} p \, dx = 0\}$. The space $H^1_0(\Omega)$ consists of functions in $H^1(\Omega)$ which vanish on the boundary in the sense of traces. To write the weak or variational formulation of the boundary value problem, we introduce the space $\mathbf{V} := [H^1_0(\Omega)]^d$ of displacements with inner product $(\cdot, \cdot)_1$ and norm $\|\cdot\|_1$ defined in the standard way; that is, $(\mathbf{u}, \mathbf{v})_1 := \sum_{i=1}^d (u_i, v_i)_1$, with the norm being induced by this inner product. A mixed formulation of the linear elastic problem is given by introducing pressure as an extra variable leading to penalized Stokes equations [4]. Defining $p := \lambda \operatorname{div} \mathbf{u}$, a mixed variational formulation of linear elastic problem is: given $\ell \in [L^2(\Omega)]^d$, find $(\mathbf{u}, p) \in \mathbf{V} \times L^2_0(\Omega)$ such that

$$\begin{aligned} A(\mathbf{u}, \mathbf{v}) + B(\mathbf{v}, p) &= \ell(\mathbf{v}), \quad \mathbf{v} \in \mathbf{V}, \\ B(\mathbf{u}, \mathbf{q}) - \frac{1}{\lambda} C(p, \mathbf{q}) &= 0, \quad \mathbf{q} \in L^2_0(\Omega), \end{aligned} \quad (13)$$

where $A(\mathbf{u}, \mathbf{v}) := 2\mu \int_{\Omega} \boldsymbol{\varepsilon}(\mathbf{u}) : \boldsymbol{\varepsilon}(\mathbf{v}) \, dx$, $B(\mathbf{v}, \mathbf{q}) := \int_{\Omega} \operatorname{div} \mathbf{v} \, \mathbf{q} \, dx$, $C(\mathbf{p}, \mathbf{q}) := \int_{\Omega} \mathbf{p} \, \mathbf{q} \, dx$ and $\ell(\mathbf{v}) := \int_{\Omega} \mathbf{f} \cdot \mathbf{v} \, dx$. Here, λ and μ are Lamé parameters, $\boldsymbol{\varepsilon}(\mathbf{u})$ is the strain of the displacement defined as $\boldsymbol{\varepsilon}(\mathbf{u}) = \frac{1}{2}(\nabla \mathbf{u} + [\nabla \mathbf{u}]^T)$ and \mathbf{f} is the prescribed body force. The existence and uniqueness of the solution of the problem (13) is shown by a standard theory of saddle point problems [3, 4].

2 Finite element discretization

We consider a quasi-uniform triangulation \mathcal{T}_h of the polygonal or polyhedral domain Ω , where \mathcal{T}_h consists of simplices, either triangles or tetrahedra. Making use of the standard linear finite element space S_h defined on the triangulation \mathcal{T}_h ,

$$S_h := \{v \in H^1(\Omega) : v|_T \in \mathcal{P}_1(T), T \in \mathcal{T}_h\},$$

and the space of bubble functions

$$B_h := \left\{ \mathbf{b}_T \in H^1(T) : \mathbf{b}_T|_{\partial T} = 0 \text{ and } \int_T \mathbf{b}_T \, dx > 0, T \in \mathcal{T}_h \right\},$$

we introduce our finite element space for the displacement as $\mathbf{V}_h = (S_h \oplus B_h)^d \cap \mathbf{V}$. The bubble function on an element T is most often defined as $\mathbf{b}_T(\mathbf{x}) = \mathbf{c}_b \prod_{i=1}^{d+1} \lambda_{T_i}(\mathbf{x})$, where $\lambda_{T_i}(\mathbf{x})$ are the barycentric coordinates of the element T associated with vertices \mathbf{x}_{T_i} of T , $i = 1, \dots, d+1$. The constant \mathbf{c}_b is computed in such a way that the value of the bubble function at the barycenter of T is one. Let N be the number of nodes in the finite element mesh, and $\{\phi_1, \dots, \phi_N\}$ be the finite element basis of S_h . Starting with the basis of S_h , we construct a dual space Q_h spanned by the basis $\{\mu_1, \dots, \mu_N\}$ so that the basis functions of S_h and Q_h satisfy biorthogonality:

$$\int_{\Omega} \mu_i \phi_j \, dx = c_j \delta_{ij}, \quad c_j \neq 0, \quad 1 \leq i, j \leq N, \quad (14)$$

where δ_{ij} is the Kronecker symbol, and \mathbf{c}_j a scaling factor, which is chosen so that $\int_{\mathcal{T}} \mu_i \, d\mathbf{x} = \int_{\mathcal{T}} \phi_i \, d\mathbf{x}$. Hence, the sets of basis functions of \mathbf{S}_h and \mathbf{Q}_h form a biorthogonal system. The basis functions of \mathbf{Q}_h is constructed locally on the reference element $\hat{\mathbf{T}}$ so that basis functions of \mathbf{S}_h and \mathbf{Q}_h have the same support [5]. We also need a subspace of \mathbf{S}_h and a subspace of \mathbf{Q}_h having zero average on Ω defined as

$$\mathbf{S}_h^0 := \left\{ \mathbf{v}_h \in \mathbf{S}_h : \int_{\Omega} \mathbf{v}_h \, d\mathbf{x} = 0 \right\}, \quad \mathbf{Q}_h^0 := \left\{ \mathbf{q}_h \in \mathbf{Q}_h : \int_{\Omega} \mathbf{q}_h \, d\mathbf{x} = 0 \right\}.$$

The first equation in (13) is discretized using a Galerkin formulation, and the second equation is discretized using a Petrov–Galerkin formulation. The Petrov–Galerkin formulation is chosen so that the pressure solution is taken from \mathbf{S}_h , whereas the test functions are taken from \mathbf{Q}_h . Hence the discrete formulation of variational equation (13) is written as: find $(\mathbf{u}_h, p_h) \in \mathbf{V}_h \times \mathbf{S}_h^0$ such that

$$\begin{aligned} \mathbf{A}(\mathbf{u}_h, \mathbf{v}_h) + \mathbf{B}_1(\mathbf{v}_h, p_h) &= \ell(\mathbf{v}_h), \quad \mathbf{v}_h \in \mathbf{V}_h, \\ \mathbf{B}_2(\mathbf{u}_h, \mathbf{q}_h) - \frac{1}{\lambda} \mathbf{C}(p_h, \mathbf{q}_h) &= 0, \quad \mathbf{q}_h \in \mathbf{Q}_h^0, \end{aligned} \tag{15}$$

where the bilinear forms $\mathbf{B}_1(\cdot, \cdot) : \mathbf{V}_h \times \mathbf{S}_h^0 \rightarrow \mathbb{R}$ and $\mathbf{B}_2(\cdot, \cdot) : \mathbf{V}_h \times \mathbf{Q}_h^0 \rightarrow \mathbb{R}$ have different domains of definition but are defined exactly as the bilinear form $\mathbf{B}(\cdot, \cdot)$ in (13). The goal of choosing the Petrov–Galerkin formulation for the pressure is to obtain a diagonal matrix corresponding to the bilinear form $\mathbf{C}(\cdot, \cdot)$ so that the degree of freedom corresponding to the pressure variable is eliminated easily.

We show the existence and uniqueness of the solution of the mixed formulation (15) using Theorem 2.

The continuity of the bilinear form $\mathbf{A}(\cdot, \cdot)$ on $\mathbf{V}_h \times \mathbf{V}_h$, of $\mathbf{B}_1(\cdot, \cdot)$ on $\mathbf{V}_h \times \mathbf{S}_h^0$, and $\mathbf{B}_2(\cdot, \cdot)$ on $\mathbf{V}_h \times \mathbf{Q}_h^0$ and of $\mathbf{C}(\cdot, \cdot)$ on $\mathbf{S}_h^0 \times \mathbf{Q}_h^0$ is straightforward. By using the Korn’s inequality, it is standard that the ellipticity of the bilinear

form $A(\cdot, \cdot)$ holds on $\mathbf{V}_h \times \mathbf{V}_h$. It remains to show that the uniform inf-sup condition holds for the bilinear form $B_1(\cdot, \cdot)$ on $\mathbf{V}_h \times \mathbf{S}_h^0$, and for the bilinear form $B_2(\cdot, \cdot)$ on $\mathbf{V}_h \times \mathbf{Q}_h^0$. As the bilinear form $B_1(\cdot, \cdot)$ on $\mathbf{V}_h \times \mathbf{S}_h^0$ satisfies the inf-sup condition uniformly with respect to the mesh size [9], we turn our attention to prove a uniform inf-sup condition for $B_2(\cdot, \cdot)$ on $\mathbf{V}_h \times \mathbf{Q}_h^0$. That means we have to show the existence of a constant $\beta > 0$ independent of the mesh-size such that

$$\sup_{\mathbf{v}_h \in \mathbf{V}_h} \frac{B_2(\mathbf{v}_h, \mathbf{q}_h)}{\|\mathbf{v}_h\|_1} \geq \beta \|\mathbf{q}_h\|_0, \quad \mathbf{q}_h \in \mathbf{Q}_h^0. \tag{16}$$

To this end, we define an operator $I_h : \mathbf{Q}_h^0 \rightarrow \mathbf{S}_h^0$ mapping every element $\mu_h = \sum_{i=1}^n c_i \mu_i$ of \mathbf{Q}_h^0 to the element $\phi_h = \sum_{i=1}^n c_i \phi_i$ of \mathbf{S}_h^0 . Using that $\int_T \phi_i \, dx = \int_T \mu_i \, dx$, $1 \leq i \leq d + 1$, we have $\int_T \phi_h \, dx = \int_T I_h \mu_h \, dx$, and hence the operator is well-defined.

For the reference triangle or reference tetrahedron, using Gauss divergence theorem, we have the following identity for any bubble function $\mathbf{b}_T \in \mathbf{B}_h$

$$\int_T \nabla \mathbf{b}_T \mu_i \, dx = (d + 2) \int_T \nabla \mathbf{b}_T \phi_i \, dx. \tag{17}$$

Using (17), we show the following properties of the operator I_h .

Lemma 3 *Given $\mathbf{v}_h \in \mathbf{V}_h$, we construct an element $\tilde{\mathbf{v}}_h \in \mathbf{V}_h$ so that there exists a constant $c > 0$ with*

$$\|\tilde{\mathbf{v}}_h\|_1 \leq c \|\mathbf{v}_h\|_1, \quad B_2(\tilde{\mathbf{v}}_h, \mathbf{q}_h) = B_2(\mathbf{v}_h, I_h \mathbf{q}_h), \quad \mathbf{q}_h \in \mathbf{Q}_h^0, \tag{18}$$

and there exist two constants $c_1 > 0$ and $c_2 > 0$ with

$$c_1 \|\mathbf{q}_h\|_0 \leq \|I_h \mathbf{q}_h\|_0 \leq c_2 \|\mathbf{q}_h\|_0. \tag{19}$$

Proof: We start with

$$B_2(\mathbf{v}_h, \mathbf{q}_h - I_h \mathbf{q}_h) = \int_{\Omega} \operatorname{div} \mathbf{v}_h (\mathbf{q}_h - I_h \mathbf{q}_h) \, dx = \sum_{T \in \mathcal{T}_h} \int_T \operatorname{div} \mathbf{v}_h (\mathbf{q}_h - I_h \mathbf{q}_h) \, dx$$

for $\mathbf{v}_h \in \mathbf{V}_h$ and $\mathbf{q}_h \in Q_h^0$. Let $\mathbf{v}_h = \mathbf{s}_h + \mathbf{b}_h$ with $\mathbf{s}_h \in S_h^d$, and $\mathbf{b}_h \in B_h^d$. Then

$$B_2(\mathbf{v}_h, \mathbf{q}_h - I_h \mathbf{q}_h) = \sum_{T \in \mathcal{T}_h} \int_T \operatorname{div} \mathbf{s}_h (\mathbf{q}_h - I_h \mathbf{q}_h) \, dx + \int_T \operatorname{div} \mathbf{b}_h (\mathbf{q}_h - I_h \mathbf{q}_h) \, dx.$$

Using that $\operatorname{div} \mathbf{s}_h$ is constant in each element T , the first integral in the right side of the above equation is zero, and hence

$$B_2(\mathbf{v}_h, \mathbf{q}_h - I_h \mathbf{q}_h) = B_2(\mathbf{b}_h, \mathbf{q}_h - I_h \mathbf{q}_h).$$

As $\mathbf{q}_h \in Q_h^0$, we write $\mathbf{q}_h = \sum_{i=1}^N q_i \mu_i$ and obtain

$$B_2(\mathbf{b}_h, \mathbf{q}_h) = \sum_{i=1}^N q_i \int_{S_i} \operatorname{div} \mathbf{b}_h \mu_i \, dx,$$

and

$$B_2(\mathbf{b}_h, I_h \mathbf{q}_h) = \sum_{i=1}^N q_i \int_{S_i} \operatorname{div} \mathbf{b}_h \phi_i \, dx,$$

where S_i denotes the support of ϕ_i or μ_i . Since \mathbf{b}_h belongs to the space of bubble functions, restricted to an element T , it is written as $\mathbf{b}_h = \mathbf{a}_T \mathbf{b}_T$ for some constant vector \mathbf{a}_T . Now we decompose the integrals inside both sums into each element

$$B_2(\mathbf{b}_h, \mathbf{q}_h) = \sum_{i=1}^N q_i \sum_{T \subset S_i} \mathbf{a}_T \cdot \int_T \nabla \mathbf{b}_T \mu_i \, dx,$$

$$B_2(\mathbf{b}_h, I_h \mathbf{q}_h) = \sum_{i=1}^N q_i \sum_{T \subset S_i} \mathbf{a}_T \cdot \int_T \nabla \mathbf{b}_T \phi_i \, dx,$$

and use the property of a bubble function (17) to get

$$B_2(\mathbf{b}_h, \mathbf{q}_h) = (d + 2)B_2(\mathbf{b}_h, I_h \mathbf{q}_h).$$

Hence defining $\tilde{\mathbf{v}}_h := \mathbf{s}_h + \frac{1}{d+2} \mathbf{b}_h$, and noting that \mathbf{s}_h and \mathbf{b}_h are linearly independent the first condition (18) is proved. The second condition (19) follows by using the fact that $\|I_h \mathbf{q}_h\|_0^2$, $\|\mathbf{q}_h\|_0^2$ and $\sum_{i=1}^N \mathbf{q}_i^2 h_i^2$ are equivalent, where h_i is the local mesh-size at the i th node of \mathcal{T}_h . ♠

A consequence of the above lemma is the following theorem.

Theorem 4 *The finite element pair (\mathbf{V}_h, Q_h^0) satisfies the inf-sup condition (16).*

Proof: Let $\mathbf{q}_h \in Q_h^0$, and $I_h \mathbf{q}_h \in S_h^0$. Since the pair (\mathbf{V}_h, S_h^0) satisfies the inf-sup condition, we can find an element $\mathbf{v}_h \in \mathbf{V}_h$ satisfying

$$B_2(\mathbf{v}_h, I_h \mathbf{q}_h) \geq c \|I_h \mathbf{q}_h\|_0^2 \quad \text{and} \quad \|\mathbf{v}_h\|_1 \leq c \|I_h \mathbf{q}_h\|_0.$$

Hence, using the properties (18) and (19) of the interpolation operator I_h , we can find an element $\tilde{\mathbf{v}}_h \in \mathbf{V}_h$ with

$$B_2(\tilde{\mathbf{v}}_h, \mathbf{q}_h) = B_2(\mathbf{v}_h, I_h \mathbf{q}_h) \geq c \|I_h \mathbf{q}_h\|_0^2 \geq c \|\mathbf{q}_h\|_0^2,$$

and

$$\|\tilde{\mathbf{v}}_h\|_1 \leq c \|\mathbf{v}_h\|_1 \leq c \|I_h \mathbf{q}_h\|_0 \leq c \|\mathbf{q}_h\|_0. \quad \spadesuit$$

Thanks to Theorems 2 and 4, the following corollary holds [6, 7, 3, 10].

Corollary 5 *The discrete problem (15) has exactly one solution $(\mathbf{u}_h, \mathbf{p}_h) \in \mathbf{V}_h \times S_h^0$, and there exists a constant c independent of Lamé parameter λ such that*

$$\|\mathbf{u}_h\|_1 + \|\mathbf{p}_h\|_0 \leq c \|\mathbf{f}\|_0.$$

Furthermore, if (\mathbf{u}, \mathbf{p}) is the solution to the problem (13), we have the following error estimate uniform with respect to λ :

$$\|\mathbf{u} - \mathbf{u}_h\|_1 + \|\mathbf{p} - \mathbf{p}_h\|_0 \leq c_1 \inf_{\mathbf{v}_h \in \mathbf{V}_h} \|\mathbf{u} - \mathbf{v}_h\|_1 + c_2 \inf_{\mathbf{q}_h \in \mathbf{S}_h^0} \|\mathbf{p} - \mathbf{q}_h\|_0,$$

where the constants c_1 and c_2 are independent of the mesh size.

Using the standard approximation properties of the spaces \mathbf{V}_h and \mathbf{S}_h^0 , we see that the approximation to the displacement converges to the exact solution with $O(h)$ in H^1 -norm.

3 Numerical results

This section illustrates the performance of the formulation discussed in the preceding sections in a numerical example [11] showing a comparison of different formulations using L^2 - and H^1 -norms. Here, the exact solution $\mathbf{u} = (\mathbf{u}_1, \mathbf{u}_2)$ is

$$\begin{aligned} \mathbf{u}_1(x, y) &:= \sin(2\pi y)(-1 + \cos(2\pi x)) + \frac{\sin(\pi x)\sin(\pi y)}{1 + \lambda}, \\ \mathbf{u}_2(x, y) &:= \sin(2\pi x)(1 - \cos(2\pi y)) + \frac{\sin(\pi x)\sin(\pi y)}{1 + \lambda} \end{aligned}$$

with $\lambda = 2499666.644443238$ and $\mu = 500.0333355557037$ (the corresponding Poisson ratio and Young's modulus are $\nu = 0.4999$ and $E = 1500$) so that a nearly incompressible response is obtained. We compute the solution by taking Ω as a unit square, where the right hand side and the Dirichlet boundary conditions are computed by using the exact solution. We have shown the discretization errors with respect to the number of elements in Figure 1. As is seen from Figure 1, the standard approach locks completely, whereas we get very good numerical approximations with our approach and mini-element. However, our formulation is more efficient as we reduce the problem to the displacement-based formulation by inverting a diagonal matrix.

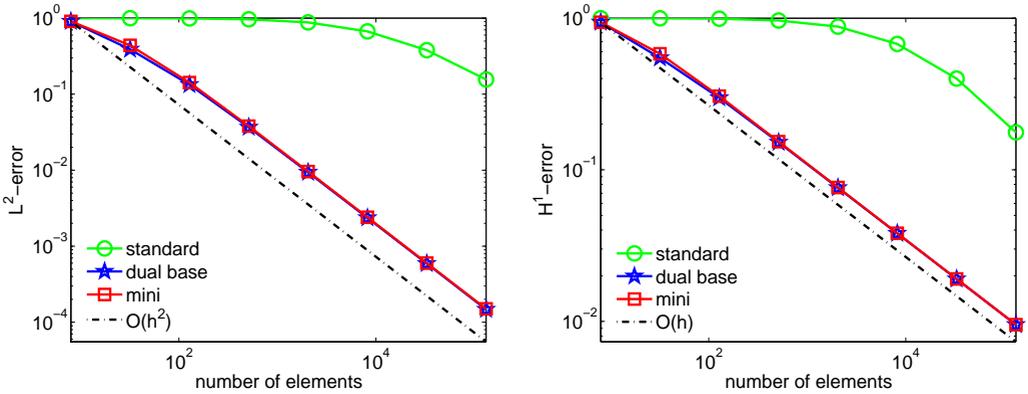


FIGURE 1: Error plot versus number of elements, L^2 -norm (left) and H^1 -norm (right).

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